

# A multi-wavelength view of the Galactic center dust ridge reveals little star formation

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## ABSTRACT

The Galactic center dust ridge consists of a narrow string of massive condensations, first identified in submillimeter dust continuum emission. To determine whether new high-mass stars are forming in this region, we performed new observations at 870  $\mu\text{m}$  with the Atacama Pathfinder EXperiment (APEX) telescope and at 8.4 GHz with the Very Large Array. We complement our data with recent maser and mid-infrared results. The ridge's clouds are dark at mid-infrared wavelengths, indicating the presence of cold, high column density material. In combination with existing temperature measurements in the dust ridge, we determine masses of the largest clouds. The results show that the dust ridge contains a very massive reservoir of molecular material. We find five radio continuum sources at 8.4 GHz in the general dust ridge vicinity, but outside of the dust ridge clouds. They are likely all excited by massive young stars, whose properties we constrain. Our observations exclude the existence of zero age main sequence stars with spectral types earlier than B0.5 within the dust ridge clouds. The only indication of ongoing high-mass star formation inside the clouds are class II methanol masers that are found in two of the clouds. Except for a weak water maser, found in previous observations, no signs of star formation are detected in the most massive cloud, M0.25+0.012.

**Key words.** Galaxy: center – stars: formation – H II regions

## 1. Introduction

The dust ridge (named by [Lis & Carlstrom 1994](#)) is an accumulation of clumpy dust condensations near the Galactic center, which are located along an arc-like narrow ridge, connecting the radio continuum sources G0.18–0.04 and Sgr B1 at Galactic longitude  $l \approx 0^\circ 5$ .

To determine the dust temperature of the molecular clouds, [Lis & Menten \(1998\)](#) and [Lis et al. \(1999\)](#) imaged the region using the Long Wavelength Spectrometer (LWS) aboard the Infrared Space Observatory (ISO) at wavelengths from 45 to 175  $\mu\text{m}$ . They found that the distribution of the emission at wavelengths longer than 70  $\mu\text{m}$  is well correlated with the submillimeter continuum emission whereas the clouds are seen in absorption against the general Galactic center background at shorter wavelengths. This implies that the clouds are much colder than and located in front of the warm dust responsible for the emission at 70  $\mu\text{m}$  and that they must have a high continuum opacity at 70  $\mu\text{m}$ . In addition, [Lis et al. \(1999\)](#) estimated the temperature at six locations by fitting the observed spectral energy distributions (SEDs) with a two-component grey-body model. They found out that the bulk of the dust has temperatures between 13 and 20 K, but that a small amount of warmer dust is also present, which potentially corresponds to emission from externally heated cloud surfaces. Ammonia observations of M0.25+0.012 by [Güsten et al. \(1981\)](#) suggest that the gaseous material has higher temperatures than the dust, probably being heated by cosmic rays.

[Molinari et al. \(2011\)](#) interpret the dust ridge as a quarter of a 100 pc size elliptical twisted ring of molecular material around

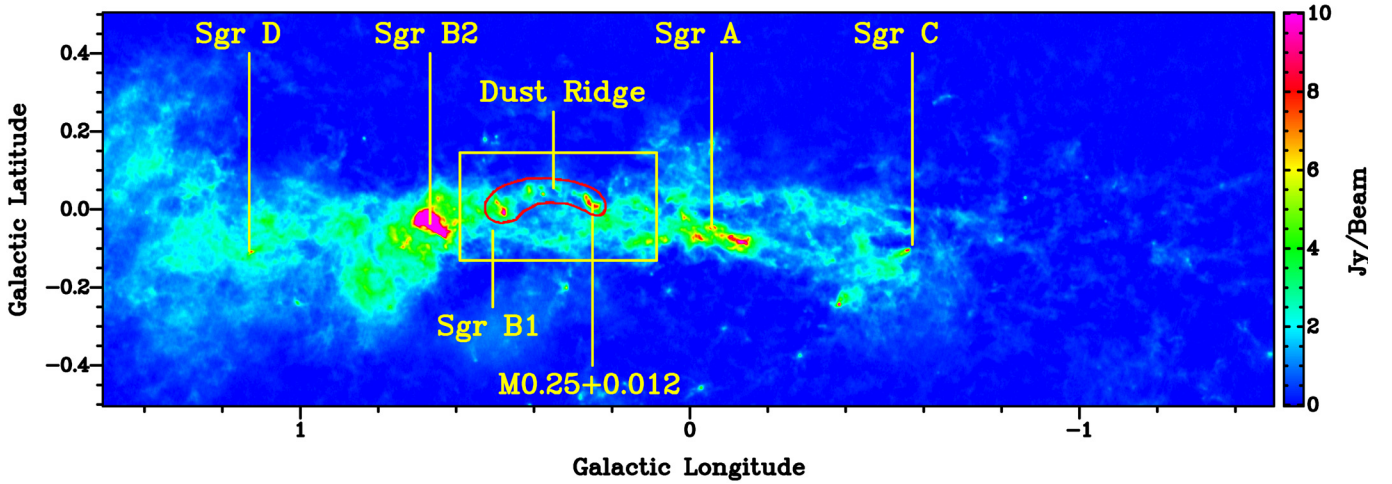
the Galactic center. Their *Herschel* observations of the central molecular zone show the dust ridge in absorption at 70  $\mu\text{m}$  and in emission at 250  $\mu\text{m}$ , confirming the low temperatures of the dust ridge clouds, estimated by [Lis et al. \(1999\)](#). The column density map of atomic hydrogen, derived from the 250  $\mu\text{m}$  map, shows strong peaks at the position of the dust ridge clouds, indicating high column densities.

[Lis & Menten \(1998\)](#) determined the distribution of the CO (2–1) and the HCO<sup>+</sup> (3–2) emission at the location of M0.25+0.012, the largest cloud in the dust ridge. The line widths of these molecules are very broad ( $\sim 30 \text{ km s}^{-1}$  FWHM), similar to typical line widths observed in giant molecular clouds (GMCs) in the Galactic center, indicating that M0.25+0.012 is likely located in the Galactic center region (i.e., at 8.5 kpc, [Ghez et al. 2008](#)).

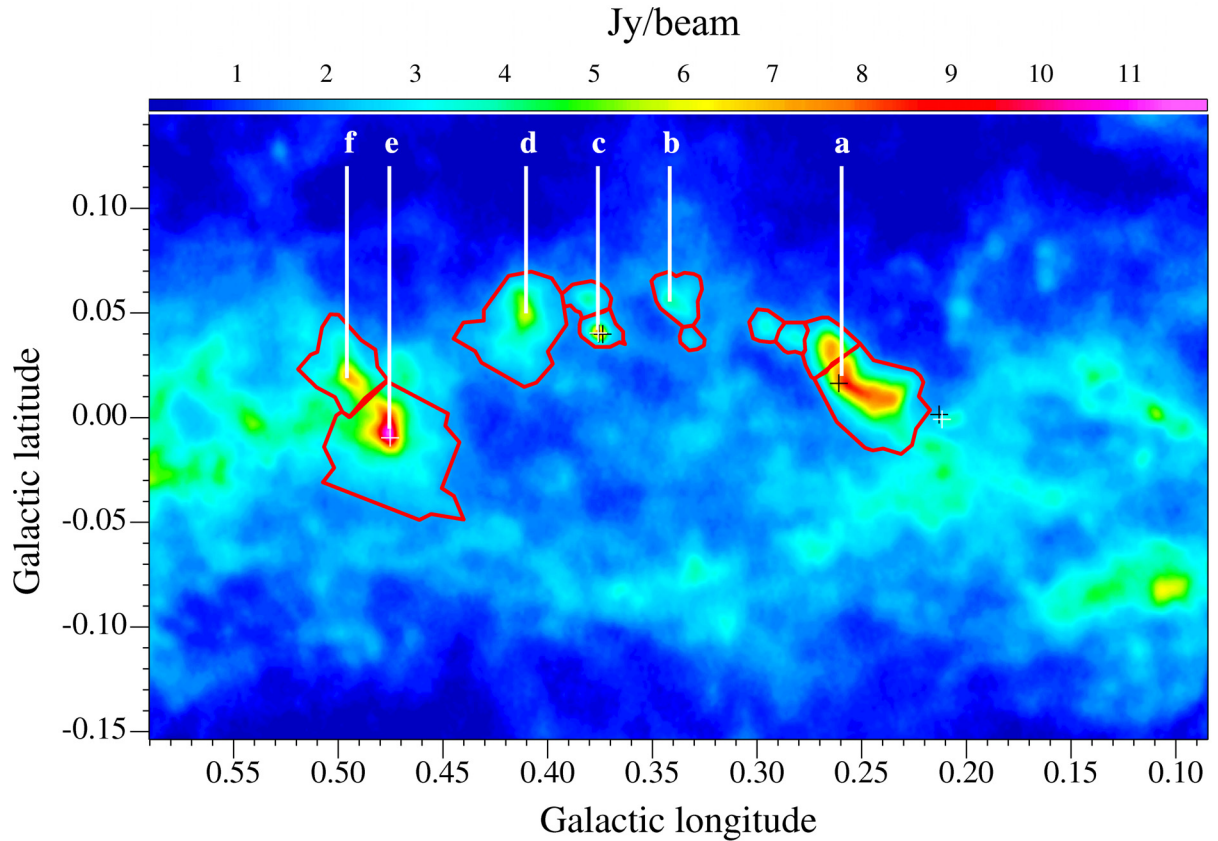
Furthermore, [Lis et al.](#) detected a large velocity gradient across the southern part of the cloud indicating streaming gas motions or several spatially overlapping velocity components. This could be a sign of a cloud collision taking place. As the cloud has not yet been warmed up by shocks, it is likely that the collision is at an early stage and may trigger substantial high-mass star formation inside the cloud in the future.

M0.25+0.012 was recently studied by [Longmore et al. \(2012\)](#), who present data from near-infrared to millimeter wavelengths. They determined global properties such as dust mass, dust temperature, and column density for this cloud and estimated the distance to the cloud from extinction measurements. Their results are consistent with the Galactic center distance of 8.5 kpc and the temperature measurements of [Lis et al. \(1999\)](#).





**Fig. 1.** 870  $\mu\text{m}$  dust continuum emission of the central molecular zone. The figure shows the position of the dust ridge (red contour) and the large molecular cloud M0.25+0.012 in relation to other famous molecular complexes: Sgr A, Sgr B1, Sgr B2, Sgr C, and Sgr D. The region within the yellow box is enlarged in Fig. 2.



**Fig. 2.** Dust continuum emission towards the dust ridge (arc of dust condensations bent to the north) at 870  $\mu\text{m}$ . The contours show the sizes of the molecular clouds over which the emission was integrated for the cloud mass determination. Water and methanol masers are marked with black and white crosses, respectively. The letters mark the positions at which Lis et al. (1999) determined the dust temperature from the SED fits.

which dust temperatures were determined from the ISO data. The lower longitude end of the ridge is located at a projected distance of around 30 pc from the Galactic center. Figure 2 shows that the 870  $\mu\text{m}$  appearance of the dust ridge region is dominated by diffuse emission with a mean value of  $\sim 1.7 \text{ Jy beam}^{-1}$ . To separate the diffuse emission from the emission of core regions and identify clumps of emission, we used the CUPID package of the starlink software with the established ClumpFind algorithm (Williams et al. 1994). We chose  $1.7 \text{ Jy beam}^{-1}$  as the lowest contour level and  $0.3 \text{ Jy beam}^{-1}$  ( $5\sigma$ ) as the gap between

contour levels. The impact of choosing different ClumpFind levels is discussed later in the text. The red contours in Fig. 2 show the boundaries of the identified clumps. Clouds a, b, and c are subdivided in several clumps. To quantify the cloud sizes, we assigned equivalent diameters based on the clump areas given by the ClumpFind algorithm.

To determine if the dust ridge clouds are gravitationally bound, the total masses of the clouds from molecular gas and dust and their virial masses have to be compared. The virial mass of a cloud is proportional to the radius of the cloud and



**Table 2.** Temperatures and masses of sub-condensations in the dust ridge in the nomenclature of [Lis et al. \(1999\)](#).

Position <sup>a</sup>	$T_{\text{Dust}}$ [K]	$F_{870\ \mu\text{m,Int.}}$ [Jy]	$M_{\text{Cloud}}$ [ $10^3 M_{\odot}$ ]	$D_{\text{Cloud}}$ ["]
a <sup>b</sup>	18	310	141	207
b	22	39	13	96
c	20	46	18	96
d	17	145	72	171
e	17	307	153	220
f	15	118	72	133

**Notes.** <sup>(a)</sup> see also Fig. 2. <sup>(b)</sup> M0.25+0.012.

the linewidth of the molecular emission squared. However, our submillimeter observations do not provide velocity information, thus prohibiting the determination of the virial mass. In a following paper, we study the kinematics of the cloud gas using data from the Millimetre Astronomy Legacy Team 90 GHz Survey (MALT90, [Foster et al. 2011](#)) and recent APEX spectral line observations to investigate whether the clouds are gravitationally bound. [Longmore et al. \(2012\)](#) showed that M0.25+0.012 (cloud a) is indeed gravitationally bound. In this paper, we can only assume that the other clouds are also gravitationally bound and not only line-of-sight superpositions. Assuming further that the dust temperature,  $T_{\text{D}}$ , does not change significantly within each cloud, we can determine the projected masses of the clouds  $M_{\text{Cloud}}$  with the formula (following [Hildebrand 1983](#))

$$M_{\text{Cloud}} = \frac{d^2 F_{870\ \mu\text{m,Int}} R}{B_{870\ \mu\text{m}}(T_{\text{D}}) \kappa_{870\ \mu\text{m}}}.$$

The integrated flux densities  $F_{870\ \mu\text{m,Int}}$  of the six clouds are measured by adding up all pixel flux density values within the boundaries of each clump. For the clouds a, b, and c, we summed up the integrated flux densities of all clumps belonging to the same cloud. For the gas-to-dust ratio,  $R$ , the opacity at  $870\ \mu\text{m}$   $\kappa_{870\ \mu\text{m}}$ , and the distance to the clouds,  $d$ , values of 100,  $0.185\ \text{m}^2\ \text{kg}^{-1}$ , and 8.5 kpc, respectively, are assumed. The results are shown in Table 2 with Cols. 2 and 3 giving the dust temperatures and the integrated flux density values at  $870\ \mu\text{m}$  of the different clouds. Column 4 contains the determined total masses of the clouds. The equivalent diameter  $D$  (Col. 5) is an estimate of the cloud size, derived from the area of the clumps given by the ClumpFind algorithm. The mass of cloud a (M0.25+0.012) is consistent with the results of [Longmore et al. \(2012\)](#).

The largest uncertainty in the mass determination arises from the determination of the cloud boundaries. Depending on the lowest contour level of the ClumpFind algorithm, the sizes of the clouds change slightly, thus including more or less diffuse emission at the boundaries of the clouds. For example, a change of the lowest contour level from  $1.7\ \text{Jy beam}^{-1}$  to  $1.8\ \text{Jy beam}^{-1}$  decreases the masses of the larger clouds by 1–5% and the smaller clouds by ~10%. The equivalent diameter decreases by 1–8% and we thus conclude that the size of the clouds is robust against small changes in the input level.

As mentioned in Sect. 2, the data reduction process was developed specifically for a better recovery of the extended emission. However, we cannot exclude that some uniform extended emission is still filtered out by the data reduction process. The masses we derive are therefore only lower limits of the real masses. Furthermore, the gas-to-dust-ratio in the central molecular zone may be lower than the standard value of 100 (determined at solar metallicity) due to the negative metallicity gradient in the disk ([Balser et al. 2011](#)). Since no metallicity

measurements exist for the dust ridge region, we adopted the commonly used value of 100 for the gas-to-dust ratio, but we note that the total masses of the clouds may be lower. We estimate the uncertainty of the total cloud masses to be of the order of ~20%.

The total mass of the dust ridge adds up to  $\sim 5 \times 10^5 M_{\odot}$ , which makes this region one of the most massive dust and gas reservoirs in the vicinity of the Galactic center. Therefore, probably an important site of the future star formation in the central molecular zone will develop in this region.

### 3.2. Maser sources in the dust ridge

[Lis et al. \(1994\)](#) conducted an observation of the cloud M0.25+0.012 at 22.2 GHz with the VLA in the B/C hybrid configuration. They found a weak water maser near the peak of the submillimeter emission with an isotropic 22 GHz line luminosity of  $1.5 \times 10^{-6} L_{\odot}$ . They deduced that this maser is most likely associated with a deeply embedded intermediate or low-mass young stellar object.

A recently conducted search for class II methanol masers at 6.7 GHz ([Caswell et al. 2010](#)) showed the existence of three methanol masers in the dust ridge region. They are located in clouds c and e, as well as in the south-west of M0.25+0.012. At the positions of the first two methanol masers also water masers have been detected ([Forster & Caswell 1999](#); [Valdettaro et al. 2001](#)). The positions of the water and methanol masers are indicated with black and white crosses, respectively, in Fig. 2.

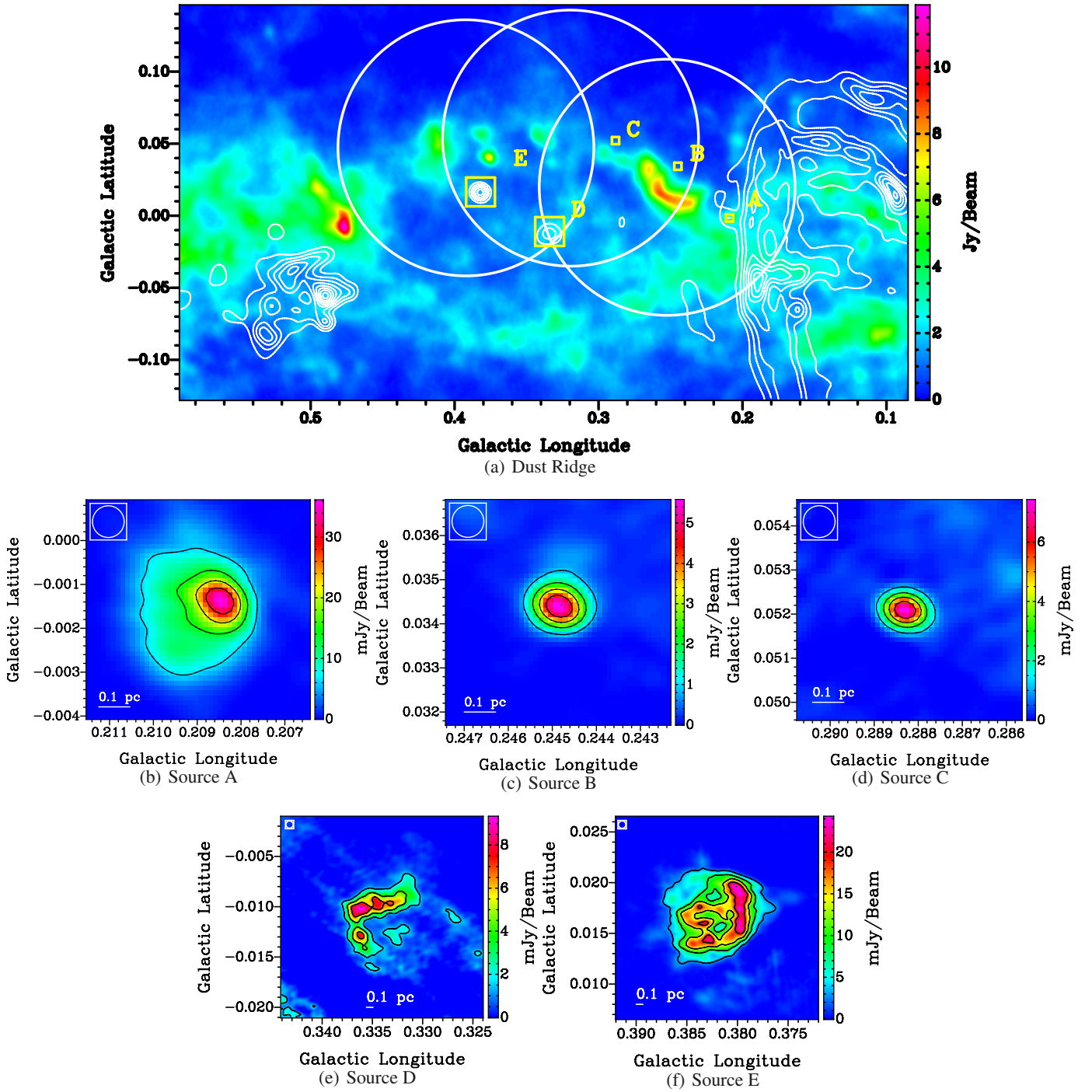
The existence of the methanol masers shows that high-mass star formation is taking place in parts of the dust ridge. However, besides the weak water maser, there were no other masers found in the most massive cloud M0.25+0.012.

### 3.3. Radio sources in the dust ridge

Figure 3 shows that the primary beams of the VLA observations cover only the dust ridge clouds a–d. Thus, we cannot make a statement about radio sources in clouds e and f from our observations. However, there are no 5 GHz or 1.4 GHz sources from the catalog of [Becker et al. \(1994\)](#) located in clouds e and f. In addition, radio emission is only detected towards the known H II region Sgr B1 in the south of clouds e and f in the 1.4 GHz map of [Yusef-Zadeh et al. \(2004\)](#) and in VLA archival data at 4.8 GHz. We, therefore, exclude the existence of H II regions in the clouds e and f.

A total of five radio sources are detected in our three 8.4 GHz VLA fields. We fitted the sources in the primary beam corrected images with Gaussian models and determined their positions and angular sizes, as well as their peak and integrated flux densities. We calculated the spectral indices of the radio sources from the integrated flux densities at 8.4 and at 5 GHz (from the catalog of [Becker et al. 1994](#)). Since [Becker et al. \(1994\)](#) list only the 1.4 GHz peak flux densities of the counterparts in their catalog, integrated flux density values at 1.4 GHz could not be included in the determination of the spectral index.

The image centered on POS1 contains three sources (labeled A, B, and C in Fig. 3) which are all offset from molecular cloud M0.25+0.012 to SWW, W and NE, respectively. Of these three sources, only the southernmost source (A) is associated with  $870\ \mu\text{m}$  dust emission, as well as water and methanol masers (see Sect. 3.2 and Fig. 2). This source is also detected at 1.4 and 5 GHz ([Becker et al. 1994](#)). Its spectral index is  $\alpha \sim +0.3$ , supporting the identification of this source as an H II region.



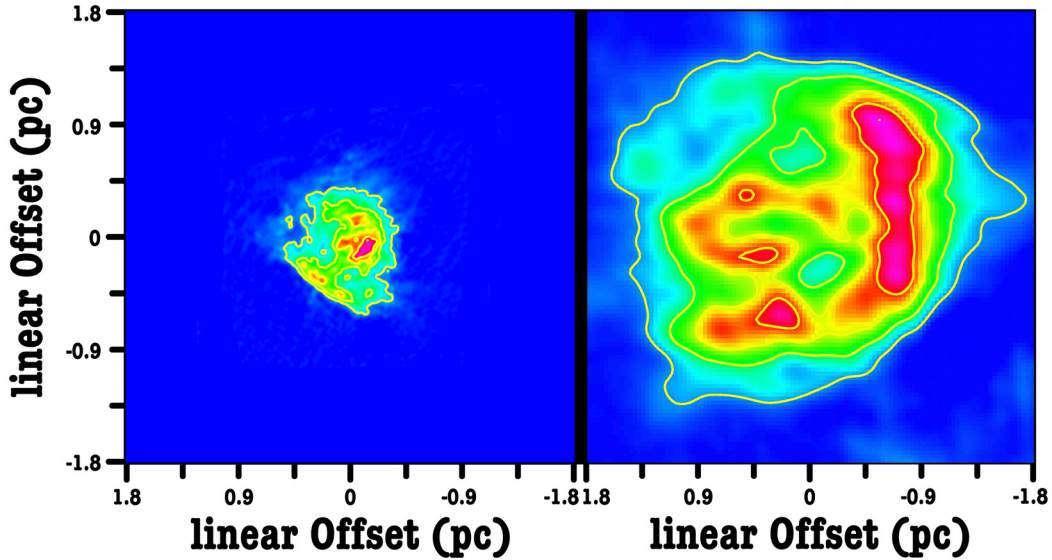
**Fig. 3.** Panel a) 870  $\mu\text{m}$  dust emission toward the dust ridge (see Fig. 2) with overlaid contours of the 1.4 GHz radio continuum emission (Yusef-Zadeh et al. 2004). Contour levels are 3% to 10%, in steps of 1%. White circles show the VLA primary beam at 8.4 GHz for the three pointings observed. Yellow boxes mark the positions and sizes of the radio continuum sources enlarged in panels b)–f). Panels b)–f) 8.4 GHz radio continuum emission toward sources A–E. Contour levels are 20% to 100%, in steps of 20% (1% corresponds to  $1.4\sigma$ ,  $0.2\sigma$ ,  $0.3\sigma$ ,  $0.6\sigma$ , and  $1.0\sigma$  in Figs. 3b–f, respectively).

The second southernmost source (B) was also detected by Lis et al. (1994) who measured a peak flux density of  $7.4 \text{ mJy beam}^{-1}$  and an integrated flux density of  $14.6 \text{ mJy}$  at 8.4 GHz. They doubted that this source has an extragalactic nature but instead assumed that they detected a compact H II region due to its relatively high flux density. The Gaussian fitting of our data results in a peak flux density of  $5.56 \pm 0.58 \text{ mJy beam}^{-1}$  and an integrated flux density of  $9.27 \pm 1.42 \text{ mJy}$ . Assuming similar uncertainties for the results of Lis et al. (1994), our results are consistent within  $3\sigma$ .

The third source detected, C, is a small, nearly circularly shaped source which is also observable in the second field POS2. Its integrated flux density is even higher than the integrated flux density of source B. Source B and C are not detected in the 1.4 GHz map of Yusef-Zadeh et al. (2004) or in the 5 GHz catalog of Becker et al. (1994). We determined lower limits for the spectral index of  $-0.2$  (source B) and  $0$  (source C), using the 5 GHz detection limit of  $9 \text{ mJy}$  as upper limit. The spectral indices suggest an interpretation of the two radio sources as H II regions.







**Fig. 4.** *Left:* the Orion nebula (M42) at 8.4 GHz, as observed with the VLA in D-array (Shepherd et al. 2001). *Right:* source E at 8.4 GHz. Contour levels are 20% to 100%, in steps of 20% in both panels (1% corresponds to  $1.4\sigma$  and  $1.0\sigma$  in the left and the right panels, respectively).

flux density of 0.98 Jy, which is comparable to the integrated flux density of source E. Therefore, if we would move the Orion nebula to the Galactic center distance of 8.5 kpc, we would see an H II region similar to source E. While the massive pre-main sequence stars in Orion are easily detectable due to their closeness to the Sun, the large distance to source E prevents the detection of the exciting stars. Thus, the only sign of the massive star formation in source E is the strong radio continuum emission. Needless to say, it would be completely impossible to detect the emission of the  $\sim 2$  million year old,  $\sim 2000$  member Orion nebula cluster at optical wavelengths, if it had been placed in the Galactic center region, given the luminosity and the pervasive high extinction toward this region.

We conclude that our data show evidence for massive star formation in a more evolved state taking place at the periphery of the dust ridge region, but outside its molecular clouds. The 8.4 GHz  $3\text{-}\sigma$  sensitivity limit of  $0.15\text{--}0.25$  mJy beam $^{-1}$  in a  $2''.6$  beam can be translated to a Lyman continuum luminosity of  $\sim 2.4\text{--}5 \times 10^{45}$  s $^{-1}$  which corresponds to a single B0.5–B1 ZAMS star. Thus, we can exclude the existence of ZAMS star with spectral types earlier than B0.5 in the dust ridge clouds.

### 3.4. Infrared emission in the dust ridge

Observations of the dust ridge at infrared wavelengths give another possibility to find newly-born stars in the clouds, since these stars heat the surrounding dust which then reemits this energy at infrared wavelengths.

Figure 5 shows the dust ridge in a GLIMPSE false-color image. The contours display the distribution of the dust emission at  $870\text{ }\mu\text{m}$ . The correspondence between the contours and the dark parts in the image in the northern part of the dust ridge is remarkable. The clouds in the southern part of the dust ridge are either not dense enough to absorb a large fraction of the diffuse infrared background or are located behind the material that emits the warm infrared background.

In addition, the infrared map shows emission at the positions of the three radio sources A, D, and E (marked with green crosses in Fig. 5). Since no maser emission is detected in the latter two sources, they might be in a more advanced stage of the star formation process.

## 4. Conclusion

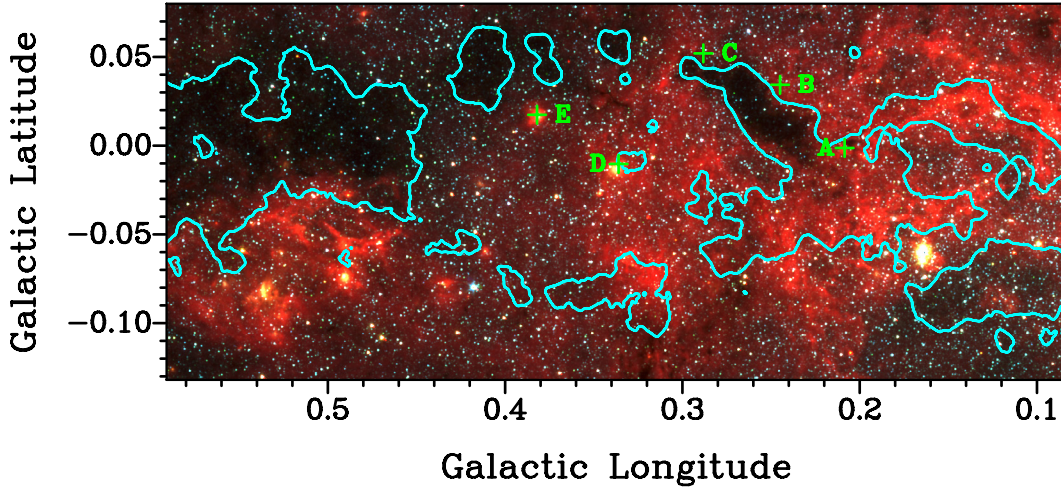
Observation of the dust ridge at  $870\text{ }\mu\text{m}$  reveal several dust clouds ordered along a narrow ridge between the radio continuum sources G0.18-0.04 and Sgr B2. The temperature of the dust condensations is very low, between 15 and 22 K, which indicates that if high-mass star formation is taking place in the clouds, it is in a very early stage. Mid-infrared observations of the dust ridge from the GLIMPSE survey show the dust ridge clouds in absorption in front of the diffuse infrared background. The agreement between the infrared dark clouds and the dust emission is remarkable.

The mass of the dust clouds are very high, ranging from  $13\,000\text{ }M_{\odot}$  to over  $150\,000\text{ }M_{\odot}$ , making this region one of the most massive reservoirs of molecular material near the Galactic center. Due to the large masses and the low temperatures, it is possible that an active site of future star formation in the central molecular zone will develop in this region.

The detection of class II methanol masers in two of the dust ridge clouds is the only sign that massive stars are born in the dust ridge. Except a weak water maser, no sign of ongoing star formation is found in the most massive dust ridge cloud, M0.25+0.012.

Observations of the dust ridge region at 8.4 GHz result in detection of five radio continuum sources, which are probably all excited by young massive stars. However, the sources are found outside the massive dust clouds, in low column density regions, indicating that the star formation at these locations is in a more evolved state than in the dense clouds. The presence of stars with spectral types earlier than B0.5 in the dust ridge clouds is excluded by our observations.

M0.25+0.012 and cloud e are by far the most massive clouds in the dust ridge. Longmore et al. (2012) identified M0.25+0.012 as a possible precursor of a young massive cluster like the Arches cluster due to its position in the radius-mass plot. The radius and mass of cloud e are comparable to the values of M0.25+0.012. Assuming that cloud e is gravitationally bound, its position in the radius-mass plot would be very close to the position of M0.25+0.012 and this cloud could therefore also form a massive young cluster in the future. While M0.25+0.012 does not show any sign of star formation, cloud e contains a methanol



**Fig. 5.** Dust ridge clouds in absorption against the infrared background (GLIMPSE RGB image: blue =  $3.6 \mu\text{m}$ , green =  $4.5 \mu\text{m}$ , red =  $8.0 \mu\text{m}$ ). The contour shows dust emission at 20% of the peak, with 1% corresponding to  $2\sigma$ . A strong correlation exists between the distribution of the dust emission and the shape of the infrared dark clouds. Green crosses and upper case letters mark the positions of the radio sources detected. Strong infrared emission is seen at the positions of the radio sources A, D, and E.

maser which indicates ongoing star formation in this cloud. We will use recent APEX spectral line observations of clouds a, d, and e in combination with data from the MALT90 survey to investigate the difference in chemistry between the clouds with and without signs of star formation and draw conclusions about the difference in the underlying physical conditions in the clouds.

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## Appendix A

Following Panagia & Walmsley (1978), the electron density  $n_e$  of a circular symmetric H II region can be determined with the formula:

$$n_e = 311.3 \times \left[ \frac{S}{\text{Jy}} \right]^{0.5} \left[ \frac{T}{10^4 \text{ K}} \right]^{0.25} \left[ \frac{D}{\text{kpc}} \right]^{-0.5} b(\nu, T)^{-0.5} \Theta_R^{-1.5} \text{ cm}^{-3}$$

with

$$b(\nu, T) = 1 + 0.3195 \log\left(\frac{T}{10^4 \text{ K}}\right) - 0.213 \log\left(\frac{\nu}{1 \text{ GHz}}\right),$$

where  $S$  denotes the flux densities,  $T$  the temperatures, and  $\Theta_R$  the angular radii of the sources.  $D$  denotes the source distance and  $\nu$  the frequency of the observation. Furthermore, the emission measure  $EM$  can be estimated from the flux density  $S$  and the angular radius  $\Theta_R$  with the following formula, which can also be obtained from Panagia & Walmsley (1978):

$$EM = 5.638 \times 10^4 \left[ \frac{S}{\text{Jy}} \right] \left[ \frac{T}{10^4 \text{ K}} \right] b(\nu, T) \Theta_R^{-2} \text{ cm}^{-6} \text{ pc}.$$

In addition, Tielens (2005) connects the emission measure with the number of Lyman continuum photons  $N_{\text{Lyc}}$  that are necessary to produce this emission measure (a constant density nebula assumed):

$$EM = 1.6 \times 10^6 \left[ \frac{n_e}{10^3 \text{ cm}^{-3}} \right]^{\frac{4}{3}} \left[ \frac{N_{\text{Lyc}}}{5 \times 10^{49} \text{ photons s}^{-1}} \right]^{\frac{1}{3}} \text{ cm}^{-6} \text{ pc}.$$

From the number of Lyman continuum photons the spectral type of the new-born star can be determined (assuming the existence of only one star in the H II region) and, together with the electron density, the mass  $M_{\text{H II}}$  of the H II region can be calculated (Tielens 2005):

$$M_{\text{H II}} \approx 80 \times \left[ \frac{n_e}{10^3 \text{ cm}^{-3}} \right]^{-1} \left[ \frac{N_{\text{Lyc}}}{5 \times 10^{49} \text{ photons s}^{-1}} \right] M_{\odot}.$$

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